ARTICLE

Neutral cumulene oxide CCCCO is accessible by one-electron oxidation of [CCCCO]^{-•} in the gas phase

Mark Fitzgerald,^{*a*} Andrew M. McAnoy,^{*a*} John H. Bowie,^{*a*} Detlef Schröder^{*b*} and Helmut Schwarz^{*b*}

^a Department of Chemistry, The University of Adelaide, South Australia, 5005. E-mail: john.bowie@adelaide.edu.au

^b Institut für Chemie, Technische Universität Berlin, D-10623, Berlin, Germany

Received 26th November 2004, Accepted 20th December 2004 First published as an Advance Article on the web 4th February 2005



Calculations at the CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory indicate that doublet [CCCCO]^{-•} is a stable species which should undergo collision-induced Franck–Condon vertical oxidation under neutralisation–reionisation conditions (⁻NR⁺) to produce both triplet CCCCO (ground state) and singlet CCCCO. Some of the neutral CCCCO species formed (particularly the triplet) should be stable for the microsecond duration of the NR experiment, whereas others will be energised (particularly the singlet) and should decompose to C₃ and CO. The [CCCCO]^{-•} radical anion has been formed in the ion source of the mass spectrometer by the reaction CH₃OCH₂C≡C–CO–CH(CH₃)₂ + O^{-•} \rightarrow [CCCCO]^{-•} + CH₃O[•] + H₂O + (CH₃)₂CH[•]. The ⁻NR⁺ spectrum of [CCCCO]^{-•} shows a recovery signal at m/z 64 corresponding to ionised CCCCO, together with a pronounced peak at m/z 36 (CCC^{+•}) produced by ionisation of CCC (formed by the reaction CCCCO \rightarrow CCC + CO). The experimental observations are in agreement with the predictions of the theoretical study.

Introduction

Cumulenes and polycarbon monoxides have been detected in circumstellar envelopes which surround red giant stars and also in dark interstellar molecular clouds.¹⁻³ Detected cumulenes include a number of C_n (n = 3 and 5), $C_n H$ (n = 2-8) and $C_n H_2$ (n = 2-4 and 6) molecules.¹⁻⁶ In contrast, there has been only limited detection of polycarbon monoxides. The linear species C_2O and C_3O [together with the associated propynal (HC=C-CHO)] have been detected towards the dark molecular cloud TMC-1.7.8 Preliminary measurements suggest that C5O may also be present, but this has yet to be confirmed.⁹ Linear C₅O has been synthesised from an anionic precursor.¹⁰ A number of theoretical studies have been devoted to polycarbon monoxides,¹¹⁻¹⁵ and the photoelectron spectra of C₂O and C₃O have also been reported.¹⁶ The rotational spectra of $C_n O$ (n = 2-9) have been detected in pyrolytic decomposition and pulsed discharge nozzle experiments.17-19

It has been suggested that the mechanism of formation of the circumstellar molecules C₃O and propynal may be linked *via* ion molecule chemistry²⁰ or by photolysis of the cluster C₃·H₂O,²¹⁻²³ and that the reaction C₃O⁺⁺ + H₂ \rightarrow C₃HO⁺ + H⁺ may provide a stellar pathway for the consumption of C₃O.²⁴ We have prepared the stable neutrals C₃O,²⁵ HCCCO²⁶ and CCCHO²⁶ in the gas phase from charged precursors by means of neutralisation reionisation mass spectrometry. It was shown that energised CCCHO may rearrange to HCCCO, which has sufficient excess energy of formation to effect dissociation to CO and C₂H,²⁶ and we have suggested that HCCCO and CCCHO may co-occur with C₃O and HC=C-CHO in dark molecular clouds.²⁶

In contrast to C₃ systems, few cumulenes containing four carbon atoms have been detected in interstellar dust clouds, circumstellar envelopes or regions of interstellar ice. The linear and rhombic isomers of C₄ have been described,²⁷ but these symmetrical species have not, as yet, been detected as stellar molecules. Neither have the cumulene oxides CCCCO or HCCCCO, although spectroscopic and/or structural details of CCCCO^{17,28-32} and HCCCCO isomers³³⁻³⁵ have been reported. The only C₄ cumulene to so far be detected in the stellar environment is HCCCC.^{5,36}

The purpose of this investigation is to determine whether the radical anion [CCCCO]⁻ can be made by an unequivocal

synthesis and if so whether this anion can be converted by oneelectron Franck–Condon vertical oxidation to a stable neutral CCCCO.

The anion [CCCCO]^{-•} has been studied previously. Theory predicts that it has a ${}^{2}\Pi$ electronic state,^{28,31} confirmed by its absorption spectrum in a neon matrix.³² It is reported that this anion may be formed by the reaction of C₄^{-•} with oxygen.³⁷ However, this synthesis does not absolutely prove the bond connectivity of the product anion and as such is not suitable for this study. Therefore we need to develop a synthesis of [CCCCO]^{-•} of confirmed bond connectivity.

Results and discussion

The aim of this study is to attempt to make the target neutral CCCCO by one-electron oxidation of the precursor anion [CCCCO]⁻⁺, *i.e.* using the neutralisation–reionisation (⁻NR⁺) procedure.³⁸⁻⁴⁰ In this procedure the anion [CCCCO]⁻⁺ is converted to neutral CCCCO by vertical one-electron oxidation. The neutral molecule is then, in turn, converted by one-electron oxidation to the cation radical [CCCCO]⁺⁺. Some of these cations will be energised and the resulting decompositions may give an insight into the structure of the first formed neutral; *i.e.* is it CCCCO, has it rearranged to another isomer, or has it decomposed to smaller neutral fragments?

This investigation is quite complex and uses experiment and theory in concert. We have chosen to report the theoretical aspects first, and then to show that the results of experiment are in accord with theory. The calculations used in the theoretical section of this paper were carried out at the CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory. Since the experimental procedure involves the sequence [CCCCO]⁻⁺ \rightarrow CCCCO \rightarrow [CCCCO]⁺⁺, the theoretical treatment follows this order.

A: The evidence based on theory

Possible rearrangements and decompositions of [CCCCO]^{-•}**.** The question that needs to be addressed first is whether [CCCCO]^{-•} is an appropriate precursor for CCCCO. In other words will [CCCCO]^{-•} survive the collisional conditions necessary to convert it to the neutral CCCCO or will it rearrange or decompose during these conditions?

| | $c^4 - c^3 - c^2 - c^1 - 0$ | $c^4 - c^3 \cdot c^2 - c_1$ | $c^{4}-c^{3}c^{2}-c^{1}$ | $c^{3}_{c^{4}}$ $c^{2}-c^{1}-0$ | $c^{3}_{c^{4}}$ c^{2} -0.01 | c^{0} | $c^{4}c^{3}-c^{2}c^{1}$ | $c^{3} < \prod_{c^{4}}^{c^{2}} c^{1} - 0$ | $o <_{C^4}^{C^3} - c^1$ |
|---|--|--------------------------------------|---------------------------------------|---|--------------------------------------|---|---|---|--------------------------------------|
| | 1. | 2 | 3 | 4 | 5 | 9 | ٦ | 8 | 6 |
| State Symmetry Energy Dipole moment | $^{2}\Pi$ $C_{\infty\nu}$ -226.94076 4.36 | C ₁ -226.77552 4.47 | $C_{s}^{2}A''$ C_{s}^{2} 2.52 | ${}^{2}B_{1}$ C_{2v} -226.90301 3.11 | C ₁ -226.79557 3.14 | ² A″ C _s -226.85894 2.97 | ${}^{2}\text{B}_{2^{\text{b}}}$ $C_{2^{\text{b}}}$ -226.78718 3.18 | C ₁ -226.88785 2.34 | C ₁ -226.78703 0.39 |
| C C C | 1.263 | | 1.267 | 1.289 | | 1.262 | 1.295 | 1.359 | 1.292 |
| 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 | 1.337 | 1.341 | | 1.496 | 1.350 | 1.367 | 1.561 | 1.049 1.576 | 1.573 |
| ひひて して ひでて | 1.281 | 1.288 | 1.267 | 1.490 1.361 1.212 | 1.558 | 1.374 | 1.295 | 1.350 | c/c.1 |
| 00000 00000 | 017:1 | 1.244 | 1.324 1.324 | 612.1 | 1.483 | 1.379 1.575 | 1.409 1.409 | 177.1 | 1.386 1.384 |
| C-C-C C-C | 180.0 | 3 021 | | 152.9 | 0 7 2 | 179.5 | 149.7 | 110.2 | 139.8 |
| $C_2 - C_2 - O$ | 180.0 | C.7/1 | | 6.70 | 0.40 | 0./61 | 149./ | 0.60 | 92.5 |
| C1-0-C2 C1-0-C2 C2 0 C3 | | 170.7 | 172.1 | | 124.5 | | | | |
| | | | 123.4 | | | | | | 94.4 |
| ひ ^っ ひって ひっして 0000 | 180.0 | 133.3 | | 180.0 | 144.6 | | 56.4 | 138.1 | |
| 0-C ⁴ -C | | | 1/2.1 | | | 55.3 | | | |
| $C^{1} - C^{2} - C^{3} - C^{4}$ | | | | 180.0 | | 0.0 | 0.0 | 0.0 | |
| | | 170.0 | 180.0 | | C 70 | | | | 6.6/1 |
| | | <i>C.C.</i> 11 | | | 77.7 | 180.0 | | | 94.4 |
| | | 180.0 | 100.0 | 0.0 | -174.6 | | 180.0 | -179.9 | |
| " CCSD(T)/aug-cc-pVDZ//B3LYF | 0/6-31G* level of theory | / including zero-poi | nt energy (B3LYP/0 | -31G*, scaled by | y 0.9804). ^{50 b} B3LY | YP/6-31G* level o | f theory. | | |

There are a number of possible $[C_4O]^{-*}$ isomers that, at least in principle, might be involved in the rearrangement of an energised $[CCCCO]^{-*}$ anion radical. These are summarised in Scheme 1 with full details listed in Table 1. Most of these are high-energy species which do not need further consideration. There are however, two species which do require further investigation, namely the cyclic anion 4^{-*} (99 kJ mol⁻¹) and the rhombic anion 8^{-*} (139 kJ mol⁻¹). The reaction profile of the interconversion $1^{-*} \rightarrow 4^{-*} \rightarrow 8^{-*}$ is shown in Fig. 1. Full geometries of the minima shown in Fig. 1 are contained in Table 1. Similar data for transition structures are listed in Table 2.

 cccco
 cccoc
 ccccc
 c c c - c - 0

 1⁻ (0 kJ mol⁻¹)
 2⁻ (434 kJ mol⁻¹)
 3⁻ (365 kJ mol⁻¹)
 4⁻ (99 kJ mol⁻¹)

5⁻ (381 kJ mol⁻¹) **6**⁻ (215 kJ mol⁻¹) **7**⁻ (403 kJ mol⁻¹) **8**⁻ (139 kJ mol⁻¹)

9^{-.} (404 kJ mol⁻¹) **10**^{-.} (not stable)

Scheme 1



Fig. 1 Rearrangement of 1^{-1} through 4^{-1} to 8^{-1} . Calculated at the CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory. Energies in kJ mol⁻¹ with respect to 1^{-1} (0 kJ mol⁻¹).

Anion 1^{-•} requires an excess energy of ≥ 143 kJ mol⁻¹ to convert through 4^{-•} to the rhomboid structure 8^{-•}, in a reaction endothermic by 139 kJ mol⁻¹. Rearrangements of negative ions of this type are the exception rather than the rule;⁴¹ those that have been confirmed were for exothermic reactions with barriers ≤ 55 kJ mol⁻¹.⁴² It thus seems likely that [CCCCO]^{-•}, once formed, should be stable under the collisional conditions required to effect the conversion of the anion to the neutral. In addition, decomposition of [CCCCO]^{-•} to [CCC]^{-•} and CO requires an excess energy of 151 kJ mol⁻¹ (Table 3). We conclude that [CCCCO]^{-•} is a suitable precursor for neutral CCCCO.

Possible rearrangements and decompositions of neutral CC-CCO. There are stable singlet and triplet forms of CCCCO,

with the triplet ground state being the more stable by 36 kJ mol^{-1} . Full details of both structures are listed in Tables 4 and 5. The results agree well with those obtained in previous studies.^{28–32} For example, the electron affinity for triplet CCCCO is calculated to be 2.86 eV [at the CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory], and this value should be compared with 2.83 eV at the CCSD/aug-cc-pVDZ level of theory.³¹ Both singlet and triplet CCCCO are very similar in structure, and are best represented by the cumulene oxide structure C=C=C=C=O.

Since the triplet and singlet structures of CCCCO are so similar in energy and geometry, it is necessary to study both the singlet and triplet potential surfaces of the C₄O system. Details of the relative energies of some triplet and singlet C₄O isomers [in comparison with that of triplet CCCCO (nominally 0 kJ mol^{-1}] are listed in Scheme 2. Full details of geometries and energies of these species are given in Tables 4 and 5 These isomers have been chosen because each, in principle, can be formed by rearrangement of energised CCCCO. Relative energies of these isomers are shown in Scheme 2.

| ссссо | сссос | ccocc | c c−c−o |
|---|--|-------------------------------|--|
| ³ 1(0 kJ mol ⁻¹) | ³ 2 (412 kJ mol ⁻¹) | ³ 3 (532 kJ mol⁻¹) | ³ 4 (not stable) |
| ¹ 1 (36 kJ mol ⁻¹) | ¹ 2 (not stable) | ¹ 3 (481 kJ mol⁻¹) | ¹ 4 (46 kJ moΓ ¹) |

С

$$c^{c}_{c} c^{-c} c^{-c}$$

Scheme 2



Fig. 2 Rearrangement of ¹1 through ¹4 to ¹8. Calculated at the CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory. Energies in kJ mol⁻¹ with respect to ¹1 (0 kJ mol⁻¹).

Triplet CCCCO has the lowest energy of the isomers listed in Scheme 2. However, the triplet forms of the other isomers are either high-energy systems or unstable at the level of theory used

| Table 2 | Energies | (E in | Hartrees)," | bond | lengths | (in | Å) ^b | and | bond |
|------------|-------------------------|--------|---------------|----------|---------|-----|-----------------|-----|------|
| angles (ir | 1 degrees) ^b | of tra | nsition state | es in Fi | g. 1 | | | | |

| | $C^{3} - C^{2} - C^{1} - O$ | $C^{3} \int_{C^{4}}^{C^{2}} C^{1} - O$ |
|-------------------------|-------------------------------------|--|
| | TS 1 ^{-•} /3 ^{-•} | TS 3 ⁻ ·/8 ⁻ · |
| State | ${}^{2}\mathbf{A}''$ | ${}^{2}\mathbf{A}''$ |
| Symmetry | $C_{\rm c}$ | C_{\circ} |
| Energy | -226.88640 | -226.88774 |
| C^1 – C^2 | 1.266 | 1.345 |
| C^1-C^4 | | 1.750 |
| $C^2 - C^3$ | 1.423 | 1.541 |
| $C^{2}-C^{4}$ | 1.970 | |
| $C^{3}-C^{4}$ | 1.313 | 1.354 |
| C^1-O | 1.224 | 1.213 |
| $C^{1}-C^{2}-C^{3}$ | 175.5 | 119.5 |
| $C^{2}-C^{3}-C^{4}$ | 92.1 | 70.4 |
| $O - C^1 - C^2$ | 175.6 | 158.5 |
| $C^1 - C^2 - C^3 - C^4$ | 0.0 | 0.0 |
| $O-C^1-C^2-C^3$ | 180.0 | 180.0 |
| | | |

^{*a*} CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory including zero-point energy (B3LYP/6-31G*, scaled by 0.9804).⁵⁰ ^{*b*} B3LYP/6-31G* level of theory.

for the calculations. Hence, triplet CCCCO should not rearrange following neutralisation from [CCCCO]⁻⁺. In contrast, the

Table 3 Relative energies (in kJ mol⁻¹) for decomposition pathways of doublet anions 1^{-1} , 4^{-1} and 8^{-1}

| $[CCCCO]^{-\bullet} \rightarrow CCC^{-\bullet} + CO$ | 151 |
|--|-----|
| $[cyc-C_3CO]^{-\bullet} \rightarrow cyc-C_3^{-\bullet} + CO$ | 95 |
| $[rhm-C_4O]^{-\bullet} \rightarrow cyc-C_3^{-\bullet} + CO$ | 56 |

CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory. Energy values were determined from the following theoretically calculated values (Hartrees): CCC⁻⁺ = -113.81418, CO = -113.06894, *cyc*-C₃⁻⁺ = -113.79772, *rhm*-C₄⁻⁺ = -151.76107, O = -74.92565.

singlet surface shows that the isomers ¹**4** and ¹**8** have relative energies of $+10 \text{ kJ mol}^{-1}$ and -24 kJ mol^{-1} with respect to singlet CCCCO (¹**1**). Can ¹**4** and ¹**8** be formed from ¹**1** during or following the one-electron oxidation of [CCCCO]^{-•}?

A simplified reaction profile of the rearrangement ${}^{1}1 \rightarrow {}^{1}4 \rightarrow {}^{1}8$ is shown in Fig. 2. Full details of ${}^{1}1$, ${}^{1}4$ and ${}^{1}8$ are given in Table 5. The geometries and energies of the transition structures are shown in Table 6. Singlet CCCCO needs an excess energy of $\geq 167 \text{ kJ mol}^{-1}$ for this reaction to proceed. If the reaction proceeds past the first transition state, ${}^{1}4$ then has sufficient excess energy to pass the second transition state and form the rhombic species ${}^{1}8$. Singlet rhombus ${}^{1}8$ is the global minimum on the singlet potential surface. We have noted previously that other rhomboid systems are also stable, *e.g.* rhombic C_4 , 27 rhombic C_4C^{43} and rhombic C_3BC .⁴⁴

The thermochemistry of possible decomposition processes of singlet and triplet CCCCO are listed in Table 7. The lowest energy decomposition pathway corresponds to loss of CO in

Table 4 Energies (E in Hartrees),^a bond lengths (in Å),^b bond angles (in degrees)^b and dipole moments (Debye)^b of triplet neutrals

| | $C^4 - C^3 - C^2 - C^1 - O$ | $C^4 - C^3 - C^2 - C^1$ | $C^4 - C^3 - C^2 - C^1$ | C^3 C^4 C^2 C^0 | $C^{3} \leq \frac{C^{2}}{C^{4}} C^{1} - O$ | $0 \stackrel{C^3}{\underset{C^4}{\overset{1}{\overset{1}{\overset{1}{}}}} C^2 - C^1$ |
|---|--|---|---|---|--|--|
| | ³ 1 | ³ 2 | ³ 3 | ³ 5 | ³ 8 | ³ 9 |
| State Symmetry Energy Dipole moment | $^{3}\Sigma$ $C_{\infty v}$ -226.83576 2.32 | ³ A" C _s -226.67866 3.78 | ³ A" C _s -226.63305 0.66 | C_1 -226.68736 0.62 | C_1 -226.76565 1.42 | ³ A" C _s -226.67452 1.22 |
| $\begin{array}{c} C^1 - C^2 \\ C^1 - C^4 \\ C^2 - C^3 \\ C^2 - C^4 \\ C^3 - C^4 \\ C^1 - O \\ C^2 - O \\ C^3 - O \\ C^4 - O \end{array}$ | 1.288 1.289 1.314 1.175 | 1.305 1.304 1.209 1.322 | 1.297 1.297 1.262 1.262 | 1.365 1.368 1.389 1.311 1.370 | 1.508 1.508 1.370 1.370 1.187 | 1.322 1.517 1.467 1.294 1.464 |
| $\begin{array}{c} C^{1}-C^{2}-C^{3}\\ C^{2}-C^{3}-C^{4}\\ C^{1}-O-C^{2}\\ C^{1}-C^{2}-O\\ C^{2}-C^{3}-O\\ C^{2}-O-C^{3}\\ C^{3}-O-C^{4}\\ O-C^{1}-C^{2}\\ O-C^{2}-C^{3}\\ O-C^{2}-C^{4}\\ O-C^{3}-C^{4}\end{array}$ | 180.0 180.0 180.0 | 171.8 166.9 1.322 | 172.8 153.5 172.8 | 59.3 120.8 151.5 | 94.1 84.5 142.3 | 133.2 100.6 85.9 |
| $\begin{array}{c} C^{1}-C^{2}-C^{3}-C^{4}\\ C^{1}-C^{2}-C^{3}-O\\ C^{1}-C^{2}-O-C^{3}\\ C^{1}-O-C^{2}-C^{3}\\ C^{2}-C^{3}-O-C^{4}\\ C^{2}-O-C^{3}-C^{4}\\ O-C^{1}-C^{2}-C^{3}\\ O-C^{2}-C^{3}-C^{4} \end{array}$ | | 180.0 180.0 | 180.0 180.0 | 0.1 180.0 | -26.9 | 180.0 0.0 |

^{*a*} CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory including zero-point energy (B3LYP/6-31G*, scaled by 0.9804).^{50 *b*} B3LYP/6-31G* level of theory.

| | 69351 | |
|--|--|--|
| | $\begin{array}{c} C_1 \\ -226, \\ 0.51 \\ 0.51 \\ 1.426 \\ 1.340 \\ 1.489 \\ 1.489 \\ 1.445 \\ 1.369 \end{array}$ | 68.4 68.4 92.2 92.2 1.4 1.4 |
| $\propto^{C_1^3}_{C^4} c^2 - c^1$ | C ₁ -226.70810 0.99 1.307 1.505 1.505 1.505 1.382 1.381 | 138.5 93.6 86.7 -136.1 -22.0 |
| $c^{3} < \frac{C^{2}}{C^{4}} c^{1} - 0$ 18 | $^{1}A_{1}$ C_{2y} -226.83125 0.69 1.468 1.468 1.427 1.182 1.182 | 114.2 66.9 147.6 0.0 180.0 evel of theory. |
| c4,c ³ -c ³ -c ² c ¹ | ¹ A ₁ C ₂₅ -226.67478 0.98 1.321 1.436 1.436 1.403 1.403 1.403 | 151.9 151.9 59.2 0.0 180.0 |
| م د ⁴ کر ³ -د ² -د ^ا ۱6 | ¹ A' C ₃ -226.75424 2.44 1.286 1.311 1.451 1.451 1.394 1.358 | 176.6 166.4 56.4 0.0 180.0 scaled by 0.9804). ⁵⁰ |
| $c_{c^{4}}^{c^{3}}$, $c^{2}-c^{1}-0$ | $^{1}A_{-1}$ $C_{2\times}$ $C_{2\times}$ 2.60 1.332 1.428 1.428 1.428 1.428 1.428 1.428 1.428 1.428 1.428 1.428 1.428 | 149.8 59.8 180.0 180.0 180.0 0.0 0.0 |
| c ⁴ →c ^{3,0} ∕c ² ~c ¹ ' 3 | ¹ A″ <i>C</i> _s <i>1</i> .37 1.293 1.293 1.293 1.288 1.288 | 170.8 126.9 170.8 180.0 180.0 180.0 ding zero-point energy |
| c ⁴ c ³ -c ² -c ¹ -0 '1 | $^{1}\Pi$ C_{∞} C_{∞} 2.12 1.290 1.292 1.312 1.176 | 180.0 180.0 180.0 G* level of theory inclu |
| | State Symmetry Energy Dipole moment $C^1 - C^2$ $C^2 - C^3$ $C^2 - C^4$ $C^2 - C^4$ | C'-C'-C' C'-C'-C' C'-C'-O C'-C'-O C'-C'-O C'-C-C' C'-C-C' O-C'-C' O-C'-C' O-C'-C' C'-C'-C' C'-C'-C' C'-C'-C' C'-C'-C' C'-C'-C' C'-C'-C' C'-C'-C' C'-C'-C' C'-C'-C' C'-C'-C' D-C'-C' C'-C'-C' D-C'-C' C'-C'-D' C'-C'-D' C'-C'-D' C'-C'-C' D-C'-C' C'-C'-C' D-C'-C' C'-C'-C' D-C'-C' D'C' D' |
| | | * CCSD(T)/; |

Table 5 Energies (*E* in Hartrees),^{*a*} bond lengths (in Å),^{*b*} bond angles (in degrees)^{*b*} and dipole moments (Debye)^{*b*} of singlet neutrals

| Table 6 | Energies | (E in | Hartrees), ^a | bond | lengths | (in) | Å) ^b | and | bond |
|------------|-------------------------|--------|-------------------------|---------|---------|-------|-----------------|-----|------|
| angles (ir | n degrees) ^b | of tra | nsition state | s in Fi | g. 2 | | | | |

| | $C^{3} - C^{2} - C^{1} - O$ | $c^{3} \underbrace{\int_{C^{4}}^{C^{2}} c^{1} - o}$ |
|---|---|---|
| | TS 11/14 | TS ¹ 4/ ¹ 8 |
| State Symmetry Energy | $^{1}A'_{C_{s}}$ -226.75845 | ¹ A' C _s -226.80881 |
| $\begin{array}{c} C^{1}-C^{2} \\ C^{1}-C^{4} \\ C^{2}-C^{3} \\ C^{2}-C^{4} \\ C^{3}-C^{4} \\ C^{1}-O \end{array}$ | 1.336 1.362 1.969 1.329 1.180 | 1.337 2.089 1.397 1.433 1.164 |
| $C^{1}-C^{2}-C^{3}$ $C^{2}-C^{3}-C^{4}$ $O-C^{1}-C^{2}$ | 149.5 94.0 158.5 | 150.2 66.5 169.5 |
| $C^1-C^2-C^3-C^4$ O-C ¹ -C ² -C ³ | 180.0 180.0 | 0.0 180.0 |

^{*a*} CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory including zero-point energy (B3LYP/6-31G*, scaled by 0.9804).⁵⁰ ^{*b*} B3LYP/6-31G* level of theory.

each case: ³1 and ¹1 require 256 and 19 kJ mol⁻¹ of excess energy in order to effect formation of CCC and CO. The singlet decomposition is a low-energy process and hence we need to know what excess energy is provided to the singlet as a consequence of the vertical one-electron oxidation. This quantity can be estimated from the difference in energy between the ground state neutral and the neutral with the anion geometry on the neutral potential surface. These calculations indicate that a singlet neutral has a Franck–Condon energy of 15 kJ mol⁻¹, an energy which is not sufficient to effect decomposition [19 kJ mol⁻¹ (Table 7)]. However, some of the singlet neutrals will experience higher internal energies due to either an initial high-energy collision or subsequent collision of the neutral. As a consequence, these neutrals may decompose to CCC and CO. In contrast, triplet neutral CCCCO should be stable under the conditions of the experiment [calculated Franck-Condon energy of 19 kJ mol⁻¹: decomposition to ³CCC + CO requires 256 kJ mol⁻¹ (Table 7)].

The conclusions from this part of this theoretical investigation are: (i) vertical oxidation of [CCCCO]^{-•} should yield some CC-CCO neutrals which are stable (mainly triplet), and some which decompose to CCC and CO (mainly singlet), (ii) rearrangement of CCCCO should not be observed, and (iii) the triplet should be stable under the experimental conditions, while the singlet will decompose to form CCC and CO provided it has an excess energy of ≥ 19 kJ mol⁻¹.

Table 7Relative energies (in kJ mol^-1) for decomposition pathways ofsinglet and triplet neutral CCCCO

| $^{1}CCCCO \rightarrow ^{1}CCC + CO$ | 19 |
|--|-----|
| $^{1}CCCCO \rightarrow ^{1}CCCO + ^{1}C$ | 500 |
| $^{1}CCCCO \rightarrow ^{3}CCCC + O$ | 691 |
| $^{3}CCCCO \rightarrow ^{3}CCC + CO$ | 256 |
| $^{3}CCCCO \rightarrow ^{1}CCCO + ^{3}C$ | 385 |
| $^{3}CCCCO \rightarrow ^{1}CCCC + O$ | 660 |

CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory. Energy values were determined from the following theoretically calculated values (Hartrees): $^{1}CCC = -113.74598$, $^{3}CCC = -113.66924$, CO = -113.06894, $^{1}CCCO = -188.92411$, $^{1}C = -37.70756$, $^{3}C = -37.76487$, $^{1}CCCC = -151.65881$, $^{3}CCCC = -151.63301$, O = -74.92565.

Possible rearrangements and decompositions of [CCCCO]^{+•}**.** The relative energies of isomeric doublet cation radicals are summarised in Scheme 3, with full details listed in Table 8. Most of the isomeric cation radicals are either unstable or of high energy with respect to $1^{+•}$ (0 kJ mol⁻¹). However, in common with data for the analogous anions and neutrals, the exceptions are $4^{+•}$ (48 kJ mol⁻¹) and $8^{+•}$ (111 kJ mol⁻¹).



9^{+.} (380 kJ mol⁻¹) **10**^{+.} (not stable)



The reaction coordinate profile for the rearrangement $1^{+*} \rightarrow 4^{+*} \rightarrow 8^{+*}$ is shown in Fig. 3. Full details of minima and transition structures shown in Fig. 3 are given in Tables 8 and 9 The rearrangement of 1^{+*} through 4^{+*} to 8^{+*} has a maximum barrier of 111 kJ mol⁻¹. The overall reaction sequence is endothermic by 111 kJ mol⁻¹. The energetics of this process are within the energy range of cation rearrangements that have been reported previously (*e.g.* ref. 27). Possible decompositions of the cation radicals shown in Fig. 3 are listed in Table 10. These are all high-energy processes and will not compete with the rearrangement process shown in Fig. 3.



Fig. 3 Rearrangement of 1^{++} through 4^{++} to 8^{++} . Calculated at the CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory. Energies in kJ mol⁻¹ with respect to 1^{++} (0 kJ mol⁻¹).

The conclusions of the theoretical study are as follows:

(1) [CCCCO]^{-•} is an appropriate precursor for CCCCO.

(2) Vertical oxidation of [CCCCO]^{-•} should yield some CCCCO neutrals which, depending on their internal energies, are either stable or decompose to CCC and CO.

(3) The [CCCCO]^{+•} cation formed from CCCCO is stable, but rearrangement to the *cyclo*-C₃ [4^{+•}] and rhomboid [8^{+•}] cations occurs below the calculated lowest energy dissociation pathways for 1^{+•}, 4^{+•} and 8^{+•}.

| | c ⁴ c ³ -c ² -c ¹ -0 1 ⁺⁺ | c ⁴ -c ³ ^c ² 0-c ¹ 2 ⁺⁺ | ^{c4} \c ³ -0-c ² . ^{c1} 3 ^{+•} | $c^{c}_{c}^{+}c^{2}-c^{1}-0$ 4+. | $c^{0}C^{3}-c^{2}C^{1}$ | c^{4} , c^{3} , c^{2} , c^{1} | $c^{3} < \int_{C^{4}}^{C^{2}} c^{1} - 0$ 8+ | $o_{c^4}^{c^3}$, c^{2-c} |
|--|---|--|--|---|---|---|--|-------------------------------------|
| State Symmetry Energy Dipole moment | $^{2}\Pi$ C_{∞} -226.49135 0.22 | C ₁ -226.37736 9.05 | C ₁ -226.25330 12.34 | ² A' C _s -226.47299 8.51 | ² A' C _s -266.37604 0.59 | ² A' <i>C</i> -226.28491 7.65 | $^{2}A_{1}$ C_{2v} -226.44899 2.86 | C ₁ -226.3466 8.66 |
| CC CC | 1.332 | | 1.335 | 1.375 | 1.336 | 1.347 | 1.511 | 1.355 |
| - - - - - - - - - - - - - - - - - - - | 1.253 | 1.227 | | 1.328 | 1.284 | 1.392 | 1.344 | 1.457 |
| - - - - - - - - - - - - - - - - - - - | 1.369 1.143 | 1.337 1.178 1.442 | 1.317 | 1.32/ 1.409 1.334 | 1.499 | 1.428 | 1.344 1.159 | 1.04.1 |
| 000 500 | | 7+++.1 | 1.261 | | 1.410 1.290 | 1.314 | | $1.360 \\ 1.360$ |
| い い い い い い い い い い い い い い い い い い い | 180.0 180.0 | 166.1 | | 158.1 67.8 | 167.0 167.3 | 162.8 150.0 | 110.4 74.3 | 141.9 |
| らっ い い い い い い い い い い い い い い い い い い い | | 157.6 | 153.8 | | | | | 9,99 |
| C2-0-C3 | | | 156.1 | | | | | 82.1 |
| 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 | 180.0 | 127.3 | | 178.7 | | 53.6 | 147.5 | |
| 0-Ct-Ct 0-Ct-Ct | | | 151.3 | | 60.2 | | | |
| C'-C ² -C' C'-C ² -C' | | | | 180.0 | 0.0 | 0.0 | 0.0 | 161.0 |
| C ¹ -C ² -O-C ³ C ¹ -O-C ² -C ³ | | 179.9 | -141.9 | | | | | |
| C ² -C ³ -0-C ⁴ C ² -0-C ⁴ | | | 8 901 | | 180.0 | | | 12.2 |
| | | -180.0 | | 180.0 | | 180.0 | 180.0 | |

Table 8 Energies (E in Hartrees), " bond lengths (in Å), " bond angles (in degrees)" and dipole moments (Debve)" of doublet cations

Table 9 Energies (*E* in Hartrees),^{*a*} bond lengths (in Å)^{*b*} and bond angles (in degrees)^{*b*} of transition states in Fig. 3

| C^{4} C^{2} C^{2} C^{1} O | $C^{3} \sqrt{\int_{C^{4}}^{C^{2}} C^{1} - O}$ |
|-------------------------------------|---|
| TS 1+•/4+• | TS 4+•/8+• |
| $^{2}A'$ | $^{2}\mathbf{A}^{\prime}$ |
| $C_{\rm s}$ | C_{s} |
| -226.45376 | -226.44905 |
| | |
| 1.335 | 1.404 |
| | 1.824 |
| 1.304 | 1.328 |
| 1.998 | |
| 1.362 | 1.362 |
| 1.149 | 1.147 |
| | |
| 170.7 | 125.2 |
| 97.1 | 76.3 |
| 175.3 | 161.3 |
| 1,010 | 10110 |
| 180.0 | 0.0 |
| 180.0 | 180.0 |
| | $\begin{array}{c} C^{3} - C^{2} - C^{1} - O \\ TS 1^{++} / 4^{++} \\ \hline \\ ^{2}A' \\ C_{s} \\ -226.45376 \\ 1.335 \\ 1.304 \\ 1.998 \\ 1.362 \\ 1.149 \\ 170.7 \\ 97.1 \\ 175.3 \\ 180.0 \\ 180.0 \\ 180.0 \\ 180.0 \\ \end{array}$ |

^{*a*} CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory including zero-point energy (B3LYP/6-31G*, scaled by 0.9804).⁵⁰ ^{*b*} B3LYP/6-31G* level of theory.

Table 10 Relative energies (in kJ mol⁻¹) for decomposition pathways of doublet cations 1^{+*} , 4^{+*} and 8^{+*}

| $[CCCCO]^{+} \rightarrow CCC^{+} + CO$ | 295 |
|--|-----|
| $[cyc-C_3CO]^{+\bullet} \rightarrow cyc-C_3^{+\bullet} + CO$ | 261 |
| $[rhm-C_4O]^{+\bullet} \rightarrow cyc-C_3^{+\bullet} + CO$ | 198 |
| $[rhm-C_4O]^{+} \rightarrow rhm-C_4^{+} + O$ | 665 |

CCSD(T)/aug-cc-pVDZ//B3LYP/6-31G* level of theory. Energy values were determined from the following theoretically calculated values (Hartrees): CCC⁺⁺ = -113.31012, CO = -113.06894, cyc-C₃⁺⁺ = -113.30453, rhm-C₄⁺⁺ = -151.27022, O = -74.92565.

B: The evidence based on experiment

The anion [CCCCO]^{-•} was made in modest yield by the process shown in eqn. (1). This is a standard process involving the reaction of O^{-•} (from N₂O)⁴⁵ with an 'activated' CH₂ group to effect loss of water. The eliminations of CH₃O[•] and (CH₃)₂CH[•] follow or accompany the expulsion of water.

$$CH_{3}O-CH_{2}-C\equiv C-CO-CH(CH_{3})_{2} + O^{-\bullet} \rightarrow CH_{3}O^{\bullet} + [CCCCO]^{-\bullet} + {}^{\bullet}CH(CH_{3})_{2} + H_{2}O$$
(1)

The collision induced (CID) negative ion mass spectrum of [CCCCO]⁻⁺ shows only a minor loss of CO, a decomposition in accordance with the bond connectivity shown and with theoretical data which predict that [CCCCO]⁻⁺ should not rearrange under CID conditions. The ⁻CR⁺ (direct conversion of [CCCCO]⁻⁺ to [CCCCO]⁺⁺) and ⁻NR⁺ (stepwise conversion of [CCCCO]⁻⁺ to CCCCO to [CCCCO]⁺⁺) spectra of [CCCCO]⁻⁺ are shown in Fig. 4.

The ${}^{-}CR^{+}$ spectrum (Fig. 4(A)) shows competitive losses of C, O, CC, CO, CCC and CCO from the molecular cation. These fragmentations are consistent with structure [CCCCO]^{+•} (1^{+•}), but some of these fragmentations (*e.g.* loss of O and CO) could also originate from energised isomers 4^{+•} or 8^{+•}. Since the anions 4^{-•} or 8^{-•} cannot be made by conventional syntheses, we are unable to compare the ${}^{-}CR^{+}$ spectra of anions 1^{-•}, 4^{-•} and 8^{-•}. The ${}^{-}NR^{+}$ spectrum (Fig. 4(B)) is weak and some minor fragment ions are difficult to see above baseline noise. Even so, the losses of C, O and CO are observed in both the ${}^{-}NR^{+}$ the ${}^{-}CR^{+}$ spectrum, a consequence of the facile loss of CO from singlet CCCCO (see Table 7). The recovery signal (*m*/*z* 64) in



Fig. 4 (A) ⁻CR⁺ spectrum of [CCCCO]⁻⁺; VG ZAB 2HF mass spectrometer. (B) ⁻NR⁺ spectrum of [CCCCO]⁻⁺; HF-ZAB/AMD 604 mass spectrometer.^{27,57}

the $^{-}NR^{+}$ spectrum is consistent with some neutral CCCCO molecules being stable for the microsecond duration of the NR experiment.

In conclusion, a combination of experiment and theory has shown that neutral CCCCO is accessible by one-electron oxidation of [CCCCO]⁻⁺. Some of the CCCCO neutrals are stable for the microsecond duration of the NR experiment, others are energised and decompose to form CCC and CO.

Theoretical and experimental details

A: Theoretical methods

Geometry optimisations were carried out with the Becke 3LYP method^{46,47} using the Dunning aug-cc-pVDZ basis set⁴⁸ within the GAUSSIAN 98 suite of programs.⁴⁹ Stationary points were characterised as minima (no imaginary frequencies). The calculated frequencies were also used to determine zeropoint vibrational energies which were used as a zero-point correction for the electronic energies. These were scaled by 0.9804.50 We have previously reported the success of the B3LYP method in predicting geometries of unsaturated chain structures, and that this method produces optimised structures, at low computational cost, that compare favourably with higher level calculations.⁵¹ More accurate energies for the B3LYP geometries were determined using the couple cluster CCSD(T) method⁵² using the Dunning aug-cc-pVDZ basis set.⁴⁸ All calculations were carried out on the Alpha Server at the Australian Partnership for Advanced Computing (APAC) National Facility (Canberra).

B: Mass spectrometric methods

A detailed description of the instrument used is given in ref. 26. In brief, the experiments were performed using a two-sector modified VG ZAB 2HF mass spectrometer with BE configuration, where B and E represent magnetic and electric sectors, respectively. The precursor anion [CCCCO]^{-•} was formed in the chemical ionisation source by the reaction between O^{-•} (from N₂O⁴⁵) and CH₃OCH₂C≡C–CO–CH(CH₃)₂ as shown in eqn. (1). CH₃OCH₂C≡C–CO–CH(CH₃)₂ was prepared by a reported procedure.⁵³

Typical source conditions were as follows: source temperature $200 \,^{\circ}$ C, repeller voltage -0.5 V, ion extraction voltage 7 kV, mass resolution $m/\Delta m \ge 1500$. Each neutral precursor was inserted into the ion source through the septum inlet, which was heated to 60 °C to give a measured pressure of ca. 10⁻⁶ Torr inside the source housing. The reagent gas $[N_2O \text{ (for } O^{-\bullet})]$ was introduced through a gas inlet into the ion source to a measured total pressure of ca. 10^{-5} Torr in the source housing. The estimated total pressure in the ion source is 10^{-1} Torr. The collisional induced dissociation (CID) spectrum was determined using B to select the parent anion, and utilising argon as the target gas in the first collision cell following B. The pressure of argon in the first cell was maintained such that 80% of the parent ion beam was transmitted through the cell. This corresponds to an average of 1.1-1.2 collisions per ion.54 Product anion peaks resulting from CID processes were recorded by scanning E. The charge reversal ($^{-}CR^{+}$) spectrum^{55,56} was recorded using single collision conditions in collision cell 1 (O2, 80% transmission of main beam).

We were unable to obtain a satisfactory neutralisation– reionisation ($^{-}NR^{+}$)³⁸⁻⁴⁰ spectrum of [CCCCO]^{-•} with the ZAB 2HF instrument (Adelaide), so this spectrum was measured using the ZAB/AMD 604 mass spectrometer (Technical University, Berlin). Full details of this instrument have been outlined previously.^{27,57} Neutralisation of anions was effected by collisional electron detachment using O₂ at 80% transmittance (of the main beam) as the collision gas in the first collision cell, while reionisation to cations was achieved by collision of the neutrals with O₂, (80% transmittance) in a second collision cell.

Acknowledgements

We thank the Australian Research Council for the award of an international link grant which enabled M. F. to spend three months at the Technical University, Berlin as an exchange postgraduate student. We also acknowledge the Australian Partnership for Advanced Computing (APAC) National Facility (Canberra) for a generous allocation of time on their Alpha Server. The Berlin group is grateful to the Deutsche Forschungsgemeinschaft and the Fonds der Chemischen Industrie for continuous support.

References

- 1 H. Olofsson, *Molecules in the Stellar Environment, Lecture Notes in Physics*, ed. O. G. Jorgenson, Springer, Heidelberg, 1994, pp. 114–133, and references cited therein.
- 2 (a) P. F. Bernath, K. H. Hinkle and J. J. Keady, *Science*, 1989, 244, 562; (b) K. H. Hinkle, *Molecules in the Stellar Environment, Lecture Notes in Physics*, ed. O. G. Jorgenson, Springer, Heidelberg, 1994, pp. 99–114, and references cited therein.
- 3 S. J. Blanksby and J. H. Bowie, Mass Spectrom. Rev., 1999, 18, 81.
- 4 A. Omont, *Molecules in the Stellar Environment, Lecture Notes in Physics*, ed. G. Jorgenson, Springer, Heidelberg, 1994, pp. 135–138, and references cited therein.
- 5 D. Smith and P. Spanel, *Mass Spectrom. Rev.*, 1995, 14, 255, and references cited therein.
- 6 For a list of stellar molecules so far identified, see: H. A. Wootten, http://www.cv.nrao.edu/~awootten/allmols.html.
- 7 M. Ohishi, H. Suzuki, S. Ishikawa, C. Yamada, H. Kanamori, W. M. Irvine, R. D. Brown, P. D. Godfrey and N. Kaifu, *Astrophys. J.*, 1991, 380, L39.

- 8 H. E. Matthews, W. M. Irvine, P. Friberg, R. D. Brown and P. D. Godfrey, *Nature*, 1984, **310**, 125.
- 9 B. E. Turner, Chemistry and Spectroscopy of Interstellar Molecules, ed. D. K. Bohme, University of Tokyo Press, Tokyo, 1992, p. 75.
- 10 S. Dua, S. J. Blanksby and J. H. Bowie, Int. J. Mass Spectrom., 2000, 195/196, 45.
- 11 R. D. Brown, F. W. Eastwood, P. S. Elmes and P. D. Godfrey, J. Am. Chem. Soc., 1983, 105, 6496.
- 12 (a) R. D. Brown, D. A. McNaughton and K. G. Dyall, *Chem. Phys.*, 1988, **119**, 89; (b) R. G. A. R. Maclagan and P. Sudkeaw, *J. Chem. Soc., Faraday Trans.*, 1993, **89**, 3325.
- 13 M. Moazzen-Ahmalsi and F. Zerbetto, J. Chem. Phys., 1995, 103, 6343.
- 14 P. Botschwina, P. Flügge and P. Sebald, J. Phys. Chem., 1995, 99, 9755.
- 15 T. Ogata, Y. Ohshima and Y. Endo, J. Am. Chem. Soc., 1995, 117, 3593.
- 16 Y. Ohshima, Y. Endo and T. Ogata, J. Chem. Phys., 1995, 102, 1493.
- 17 J. M. Oakes and G. B. Ellison, *Tetrahedron*, 1986, 42, 6263.
- 18 R. D. Brown, P. D. Godfrey, P. S. Elmes, M. Podler and L. M. Tack, J. Am. Chem. Soc., 1985, 107, 4112.
- 19 D. W. Ewing, J. Am. Chem. Soc., 1989, 111, 8809.
- 20 (a) S. Petrie, R. P. A. Bettens, C. G. Freeman and M. J. McEwan, *Mon. Not. R. Astron. Soc.*, 1993, **264**, 862; (b) S. Petrie, R. P. A. Bettens, C. G. Freeman and M. J. McEwan, *J. Phys. Chem.*, 1993, **97**, 13673.
- 21 S. Petrie, R. P. A. Bettens, C. G. Freeman and M. J. McEwan, J. Phys. Chem., 1993, 97, 13673.
- 22 R. G. A. Maclagan, M. J. McEwan and G. B. Scott, *Chem. Phys. Lett.*, 1995, 240, 185.
- 23 S. Ekern, J. Szczepanski and M. Vale, J. Phys. Chem., 1996, 100, 16109.
- 24 B. E. Turner and L. M. Zuirys, Interstellar Molecules and Astrochemistry, in *Galactic and Extragalactic Radio Astronomy*, ed. G. L. Verschuur and K. I. Kellerman, Springer-Verlag, Berlin, 1988.
- 25 S. J. Blanksby, S. Dua and J. H. Bowie, *Rapid Commun. Mass Spectrom.*, 1999, 13, 2249.
- 26 S. Peppe, S. J. Blanksby, S. Dua and J. H. Bowie, J. Phys. Chem. A, 2000, 104, 5817.
- 27 S. J. Blanksby, D. Schröder, S. Dua, J. H. Bowie and H. Schwarz, J. Am. Chem. Soc., 2000, 122, 7105.
- 28 H. Kannari, K. Aoki, K. Hashimoto and S. Ikuta, *Chem. Phys. Lett.*, 1994, **222**, 313.
- 29 N. Moazzen-Ahmadi and N. Zerbetto, J. Chem. Phys., 1995, 103, 6343.
- 30 S. T. Brown, J. C. Rienstra-Kiracofe and H. F. Schaefer III, J. Phys. Chem., 1999, 103, 4065.
- 31 J. C. Rienstra-Kiracofe, G. B. Ellison, B. C. Hoffman and H. F. Schaefer III, J. Phys. Chem., 2000, 104, 2273.
- 32 E. Riaplov, M. Wyss, N. M. Lakin and J. P. Maier, J. Phys. Chem., 2001, 105, 4894.
- 33 H. Kohguchi, Y. Ohshima and Y. Endo, J. Chem. Phys., 1994, 101, 6463.
- 34 Y. Hidaka, Y. Henmi, T. Ohonishi, T. Okuno and T. Koike, *Combust. Flame*, 2002, **130**, 62.
- 35 M. Fitzgerald, S. Dua and J. H. Bowie, Int. J. Mass Spectrom., 2003, 228, 467.
- 36 J. H. Bieging and M. Tafella, *Astron. J.*, 1993, **105**, 576, and references cited therein.
- 37 N. Watanabe, H. Shiromara, Y. Negishi, Y. Achiba, N. Kobayashi and Y. Kaneko, Z. Phys. D. Suppl., 1993, 87, 6934.
- 38 C. Wesdemiotis and F. W. McLafferty, Chem. Rev., 1987, 87, 485.
- 39 (a) D. V. Zagorevskii and J. L. Holmes, *Mass Spectrom. Rev.*, 1994, 13, 133; (b) N. Goldberg and H. Schwarz, *Acc. Chem. Res.*, 1994, 27, 347.
- 40 (a) For NR nomenclature, see: C. A. Schalley, G. Hornung, D. Schröder and H. Schwarz, *Int. J. Mass Spectrom. Ion. Process.*, 1998, 172, 181; (b) D. V. Zagorevskii and J. L. Holmes, *Mass Spectrom. Rev.*, 1999, 18, 87.
- 41 S. Dua, S. J. Blanksby, S. Peppe, A. M. McAnoy and J. H. Bowie, *Curr. Org. Chem.*, 2003, 7, 1545.
- 42 S. Dua and J. H. Bowie, J. Chem. Soc., Perkin Trans. 2, 2001, 827.
- 43 S. Dua and J. H. Bowie, J. Phys. Chem. A., 2002, 106, 1374.
- 44 A. M. McAnoy, J. H. Bowie and S. J. Blanksby, J. Phys. Chem. A, 2003, 107, 10149.
- 45 J. H. Dawson and K. R. Jennings, J. Chem. Soc. Faraday Trans. 2, 1976, 72, 700.
- 46 A. D. Becke, J. Chem. Phys., 1993, 98, 5648.
- 47 P. J. Stevens, F. J. Devlin, C. F. Chablowski and M. J. Frische, J. Phys. Chem., 1994, 98, 11623.

- 48 (a) T. H. Dunning, J. Chem. Phys., 1989, 90, 1007; (b) D. E. Woon and T. H. Dunning, J. Chem. Phys., 1993, 98, 1358; (c) T. H. Dunning, K. A. Peterson and D. E. Woon, Basis Sets: Correlation Consistent, in Encyclopedia of Computational Chemistry, ed. P. v. R. Schleyer, Wiley: Chichester, 1998.
- 49 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, V. G. Zakrzewski, J. A. Montgomery, Jr., R. E. Stratmann, J. C. Burant, S. Dapprich, J. M. Millam, A. D. Daniels, K. N. Kudin, M. C. Strain, O. Farkas, J. Tomasi, V. Barone, M. Cossi, R. Cammi, B. Mennucci, C. Pomelli, C. Adamo, S. Clifford, J. Ochterski, G. A. Petersson, P. Y. Ayala, Q. Cui, K. Morokuma, D. K. Malick, A. D. Rabuck, K. Raghavachari, J. B. Foresman, J. Cioslowski, J. V. Ortiz, A. G. Baboul, B. B. Stefanov, G. Liu, A. Liashenko, P. Piskorz, I. Komaromi, R. Gomperts, R. L. Martin, D. J. Fox, T. Keith, M. A. Al-Laham, C. Y. Peng, A. Nanayakkara, C. Gonzalez, M. Challacombe, P. M. W. Gill, B. G. Johnson, W. Chen,

M. W. Wong, J. L. Andres, M. Head-Gordon, E. S. Replogle and J. A. Pople, *GAUSSIAN 98 (Revision A.9)*, Gaussian, Inc., Pittsburgh, PA, 1998.

- 50 M. W. Wong, Chem. Phys. Lett., 1996, 256, 391.
- 51 S. J. Blanksby, S. Dua and J. H. Bowie, J. Phys. Chem., 1999, 103, 5161.
- 52 M. J. O. Deegan and P. J. Knowles, *Chem. Phys. Lett.*, 1994, **1227**, 321, and references cited therein.
- 53 C. Chowhury and N. G. Kunda, Tetrahedron Lett., 1996, 37, 7323.
- 54 J. L. Holmes, Org. Mass Spectrom., 1985, 20, 69.
- 55 (a) J. H. Bowie and T. Blumenthal, J. Am. Chem. Soc., 1975, 97, 2959; (b) J. E. Szulejko, J. H. Bowie, I. Howe and J. H. Beynon, Int. J. Mass Spectrom. Ion Phys., 1980, 13, 76.
- 56 M. M. Bursey, Mass Spectrom. Rev., 1990, 9, 555.
- 57 C. A. Schalley, G. Hornung, D. Schröder and H. Schwarz, *Chem. Soc. Rev.*, 1998, 27, 91.